



# Multi-loop unitarity via computational algebraic geometry

Loopfest XIII, NYC, 2014

Yang Zhang

Niels Bohr Institute

Based on

(2-loop 4-point) arXiv:1202.2019, Simon Badger, Hjalte Frellesvig and YZ

(algebraic geometry methods) arXiv:1205.5707, YZ

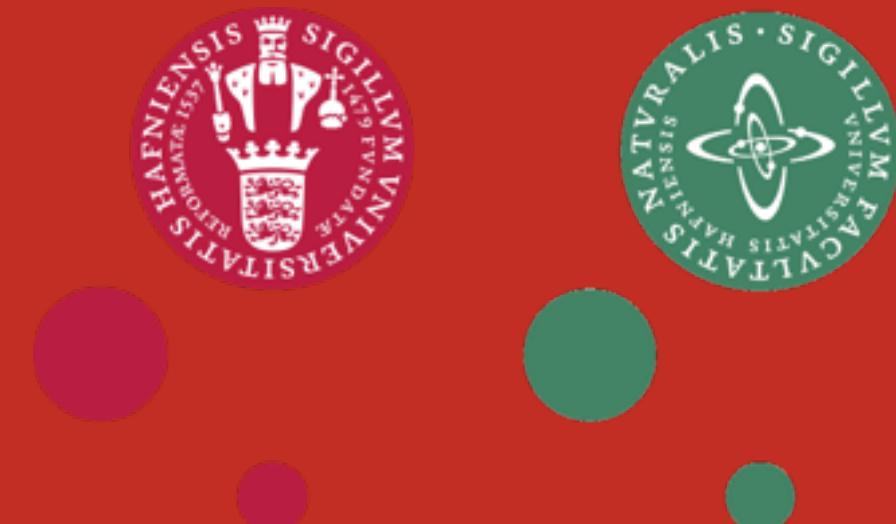
(3-loop 4-point) arXiv:1207.2976, Simon Badger, Hjalte Frellesvig and YZ

(global structure) arXiv:1302.1023, Rijun Huang and YZ

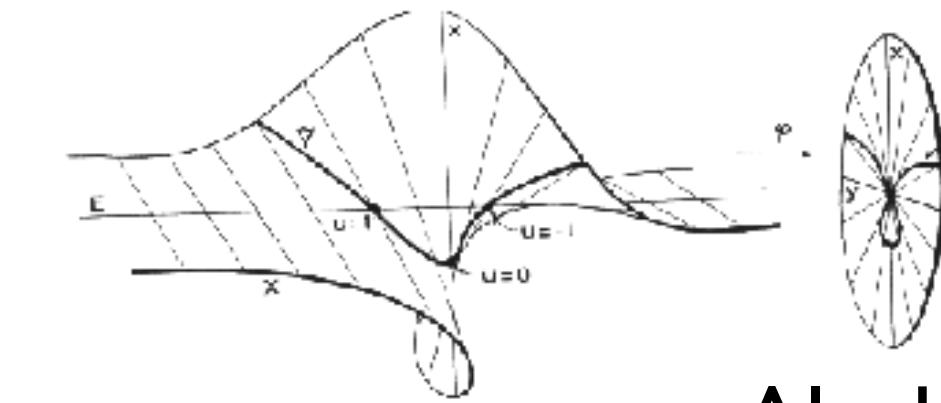
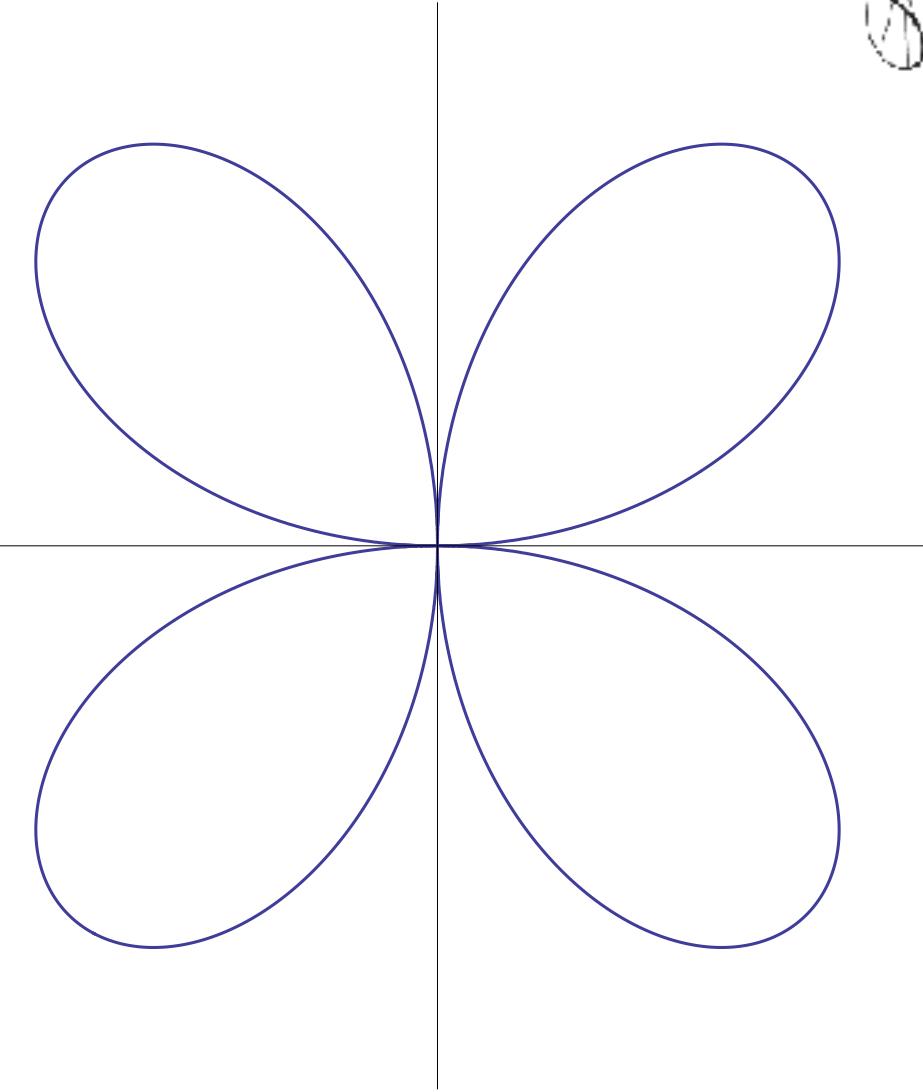
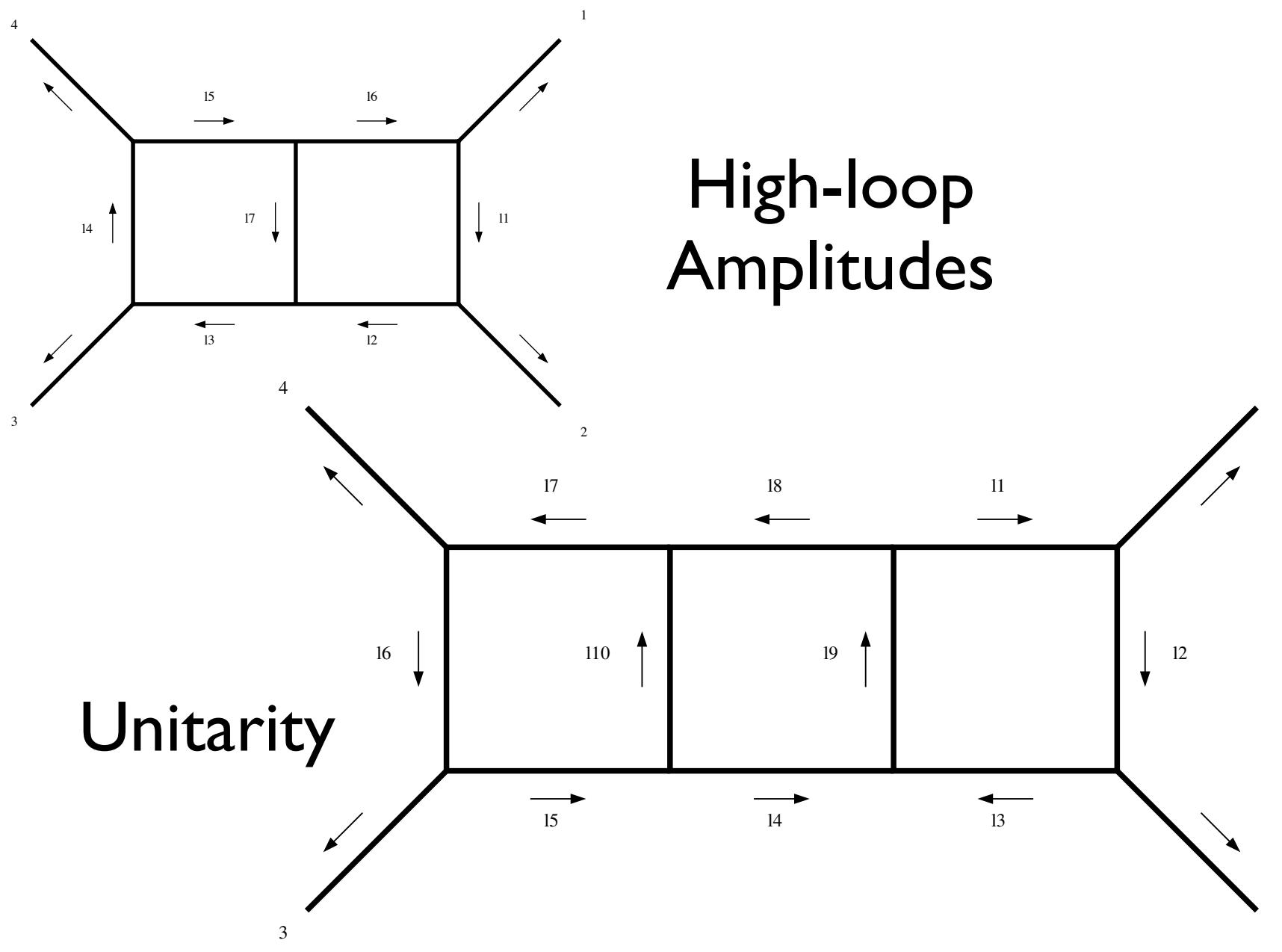
(2-loop 5-point QCD) arXiv:1310.1051, Simon Badger, Hjalte Frellesvig and YZ

(maximal cut) arXiv:1310.6006, Mads Sogaard and YZ

and works under progress...



# Outline



Unitarity

High-loop  
Amplitudes

Algebraic  
geometry

Gröbner Basis  
Primary Decomposition  
Affine Variety Structure  
Multivariate residue

- Integrand reduction at one loop, review
- Integrand reduction at n loop by algebraic geometry
- Examples: 2-loop 5-gluon planar QCD, 3-loop 4-point triple-box ...

# Unitarity at one-loop

$D = 4$

$$A^{(1)} = c_{\text{box}} \cdot \text{[Diagram: four-point box with internal lines]} + c_{\text{tri}} \cdot \text{[Diagram: three-point triangle with internal line]} + c_{\text{bub}} \cdot \text{[Diagram: two-point loop with internal line]} + \dots$$

- no pentagon, hexagon ...
- **scalar** integral (numerator is one.)

Unitarity:

Determine ‘c’ coefficients  
from **on-shell** cut solutions  
and **tree amplitudes**

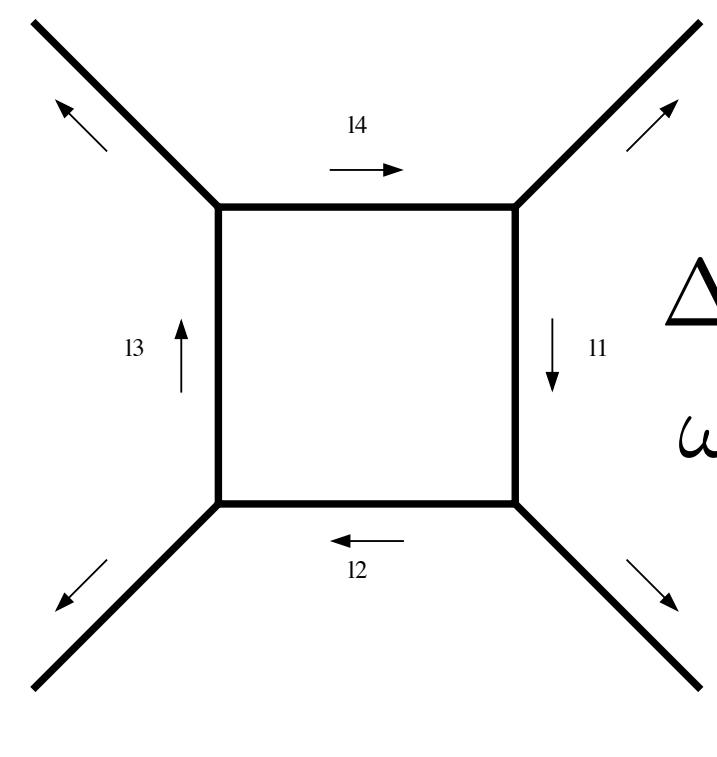
quadruple cut  $\rightarrow c_{\text{box}}$

triple cut  $\rightarrow c_{\text{tri}}$

double cut  $\rightarrow c_{\text{bub}}$

# Integrand reduction: box

Integrand-level reduction, Ossola, Papadopoulos and Pittau (OPP), 2006  
 Giele, Kunszt, Melnikov, 2008

$$A^{(1)} = \int \frac{d^4 k}{(2\pi)^4} \frac{N(k)}{D_1 D_2 D_3 D_4}$$


$$\begin{aligned} N(k) &= \Delta_{1234}(k) + \sum_{i_1 < i_2 < i_3} \Delta_{i_1 i_2 i_3}(k) \prod_{i \neq i_1, i_2, i_3} D_i + \sum_{i_1 < i_2} \Delta_{i_1 i_2}(k) \prod_{i \neq i_1, i_2} D_i \\ &= \Delta_{1234}(k) + O(D_1, D_2, D_3, D_4) \end{aligned}$$

$\Delta_{1234}(k)$  is a polynomial in scalar products (SP).  $\mathbb{SP} = \{k \cdot P_1, k \cdot P_2, k \cdot P_3, k \cdot \omega\}$

$\omega$  is auxiliary,  $(\omega \cdot P_i) = 0, i = 1, 2, 3, 4$

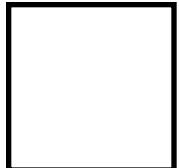
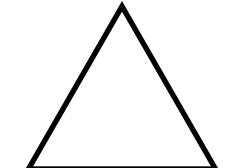
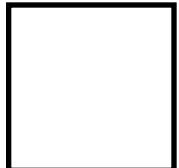
$2(k \cdot P_1) = D_4 - D_1 - P_1^2$	<hr style="border: 1px solid red;"/>	$2(k \cdot P_2) = D_1 - D_2 + P_2^2$
$2(k \cdot P_3) = D_2 - D_3 + 2P_2 \cdot P_3 + P_3^2$	<hr style="border: 1px solid red;"/>	$\text{reducible scalar product (RSP)}$ $\text{irreducible scalar product (ISP)}$

$$\Delta_{1234}(k) = \sum_i c_i (k \cdot \omega)^i \quad \Delta_{1234}(k) \text{ is a polynomial in ISP only.}$$

integrand basis

$$\boxed{\Delta_{1234}(k) = c_0 + c_1(k \cdot \omega)}$$

# One loop, other diagrams

Dimension	Diagram	# SP (ISP+RSP)	#terms in integrand basis (non-spurious + spurious)	# Solutions (dimension)
4		4 (1+3)	2 (1+1)	2 (0)
4		4 (2+2)	7 (1+6)	1 (1)
4		4 (3+1)	9 (1+8)	1 (2)
4-2 $\epsilon$		5 (2+3)	5 (3+2)	1 (1)

- straightforward to obtain **integrand basis, unitarity cut** solutions
- all one-loop **master integrals** are known
- **c coefficients** can be automatically computed by public codes

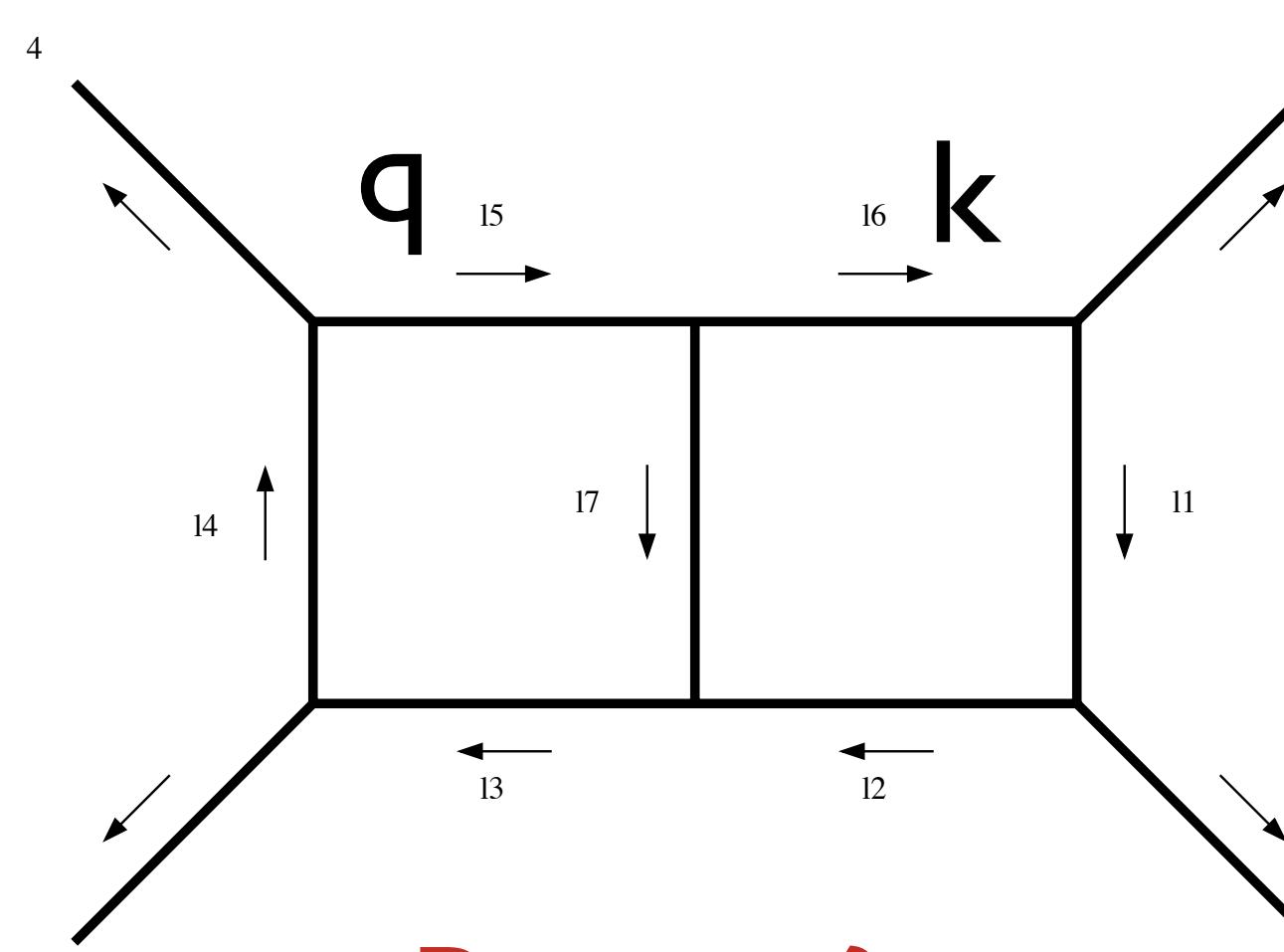
- ‘NGluon’, Badger, Biedermann, and Uwer
- ‘CutTools’, Ossola, Papadopoulos, and Pittau
- ‘GoSam’, Cullen, Greiner, Heinrich, Luisoni, and Mastrolia
- ...

Generalization to  
higher loops?

# Example: 4D massless two-loop hepta cut

P. Mastrolia, G. Ossola, 2011

S. Badger, H. Frellesvig, YZ, 2012



**Basis =?**

7 cut-equations in 8 SP's

$$\text{SP} = \{k \cdot P_1, k \cdot P_2, k \cdot P_4, k \cdot \omega, q \cdot P_1, q \cdot P_2, q \cdot P_4, q \cdot \omega\}$$

4 cut-equations to identify 4 RSP's

4 ISP's

$$\text{ISP} = \{k \cdot P_4, k \cdot \omega, q \cdot P_1, q \cdot \omega\}$$

3 cut-equations for ISP's

$$(k \cdot \omega)^2 = (k \cdot P_4 - t/2)^2 \quad (1)$$

$$^2 (q \cdot \omega)^2 = (q \cdot P_1 - t/2)^2 \quad (2)$$

$$(k \cdot \omega)(q \cdot \omega) = -\frac{t^2}{4} + \frac{t(k \cdot P_4)}{2} + \frac{t(q \cdot P_1)}{2} + \left(1 + \frac{2t}{s}\right)(k \cdot P_4)(q \cdot P_1) \quad (3)$$

Naive guessing: all renormalizable monomials which do **NOT** contain  $(k \cdot \omega)^2$ ,  $(q \cdot \omega)^2$  or  $(k \cdot \omega)(q \cdot \omega)$ .

$$\Delta_{\text{dbox}} = (k \cdot P_4)^m (q \cdot P_1)^n (k \cdot \omega)^\alpha (q \cdot \omega)^\beta$$

$$m + \alpha \leq 4, n + \beta \leq 4, m + n + \alpha + \beta \leq 6$$

$$(\alpha, \beta) = (0, 0), (1, 0), (0, 1)$$

**56 terms? wrong...**

# Example: 4D massless two-loop hepta cut

S. Badger, H. Frellesvig, YZ, 2012

3 cut-equations for ISP's, and their combinations

$$(k \cdot \omega)^2 = (k \cdot P_4 - t/2)^2 \quad (1)$$

$$(q \cdot \omega)^2 = (q \cdot P_1 - t/2)^2 \quad (2)$$

$$(k \cdot \omega)(q \cdot \omega) = -\frac{t^2}{4} + \frac{t(k \cdot P_4)}{2} + \frac{t(q \cdot P_1)}{2} + \left(1 + \frac{2t}{s}\right)(k \cdot P_4)(q \cdot P_1) \quad (3)$$

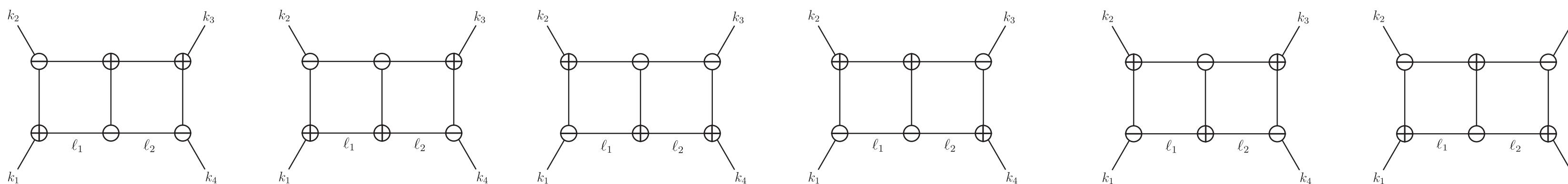
reduced

$$\boxed{(1) \times (2) - (3)^2}$$

$$4(k \cdot P_4)^2(q \cdot P_1)^2 = -2s(k \cdot P_4)^2(q \cdot P_1) - 2s(k \cdot P_4)(q \cdot P_1)^2 - st(k \cdot P_4)(q \cdot P_1)$$

We have to “exhaust” all combinations...

Finally, we determine that the basis contains 32 terms



6 families of hepta-cut solutions, Laurant series contains 38 terms

Solving 38 linear equations for 32 coefficients, done!

Messy, not automatic!

# Gröbner basis and integrand basis

arXiv:1205.5707, YZ

arXiv:1205.7087, Mastrolia, Mirabella, Ossola and Peraro

## Synthetic polynomial division

$$I = \langle D_1, \dots, D_k \rangle = \left\{ \sum_{i=1}^k g_i D_i \mid \forall g_i \in R \right\}$$

$$\int \frac{d^4 k}{(2\pi)^4} \int \frac{d^4 q}{(2\pi)^4} \frac{N}{D_1 D_2 \dots D_7}, \quad N = Q + \Delta_{\text{dbox}}, \quad Q \in I$$

$N$  divided by  $\{D_1, \dots, D_k\}$ :

Define a **monomial order**, and recursively perform  $N/D_1, \dots, N/D_k$ . Finally,  **Euclidean division**  
the division process will stop and we have

$$N = f_1 D_1 + \dots + f_k D_k + r'$$

where  $r'$  is the **remainder**.  $\Delta_{\text{dbox}} = r' ???$

In most cases, it does not work since it stops too early,  
unless we are using Gröbner basis.

$$I = \langle D_1, \dots, D_k \rangle = \langle g_1, \dots, g_m \rangle$$

$$N = q_1 g_1 + \dots + q_m g_m + r$$

- $r$  is uniquely determined.

$$\boxed{\Delta_{\text{dbox}} = r}$$

**Gröbner basis**  
**'good' generators**

$$(y^3 \quad x - 2y^2) = (x^3 - 2xy \quad x^2y - 2y^2 + x) \begin{pmatrix} -\frac{1}{4} & \frac{1}{4}xy & -\frac{1}{2}y^3 & y^2 \\ \frac{1}{4}x^2 & -\frac{1}{2}y + \frac{1}{2}xy^2 & 1 - xy \end{pmatrix}$$

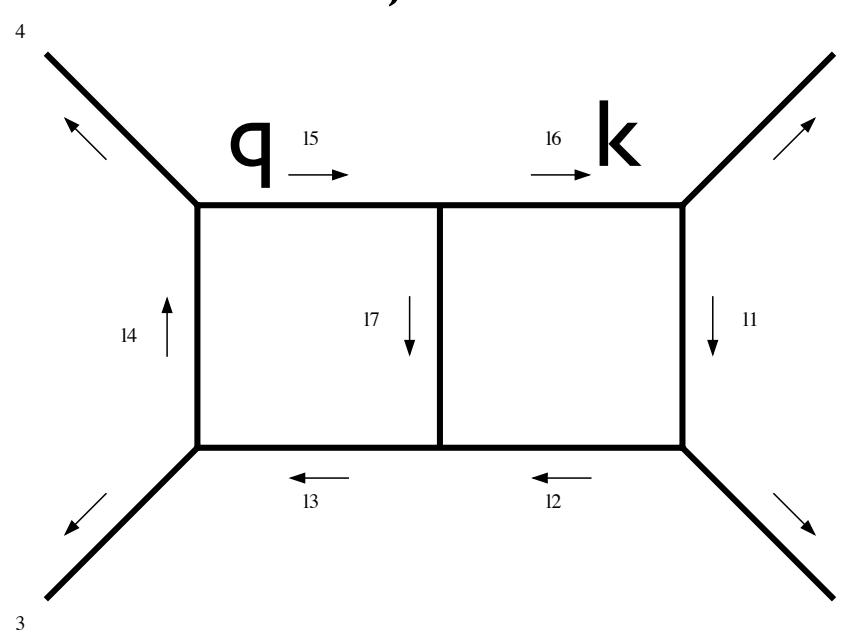
- If  $N \in I$ ,  $r = 0$ .

Toy Model:  $N = xy^3$ ,  $I = \langle x^3 - 2xy, x^2y - 2y^2 + x \rangle$ . Direct synthetic division of  $N$  towards  $\{x^3 - 2xy, x^2y - 2y^2 + x\}$  gives  $\textcolor{red}{r'} = \textcolor{red}{xy^3}$ .

But the Gröbner basis is  $I = \langle y^3, x - 2y^2 \rangle$ , and the synthetic division of  $N$  on Gröbner basis gives  $\textcolor{red}{r} = 0$ . So  $N \in I$ .

# Grobner basis: dbox example

arXiv:1205.5707, YZ



4 ISP's     $\text{ISP} = \{k \cdot P_4, k \cdot \omega, q \cdot P_1, q \cdot \omega\}$

$$N = q_1 g_1 + \dots q_k g_k + \Delta_{\text{dbox}}$$

$N$  contains 160 terms where  $\Delta_{\text{dbox}}$  contains 32 terms.

In principle, it works for arbitrary number of loops, any dimension.  
Automated by the package: '**BasisDet**'

<http://www.nbi.dk/~zhang/BasisDet.html>, YZ 2012

Dimension  
propagators,  
kinematics



Integrand  
basis

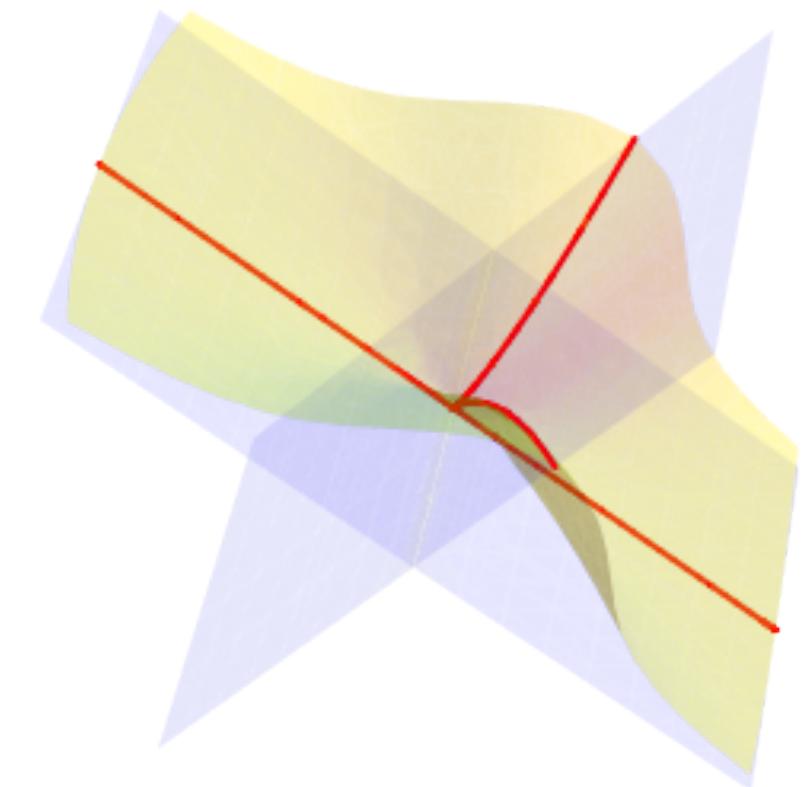
Can also find ISP  
automatically!

# Primary decomposition

arXiv:1205.5707, YZ

Find the number of branches of unitarity solutions

$I = \langle x^2 - y^2, x^3 + y^3 - z^2 \rangle$ . How many (irreducible) curves are there in  $\mathcal{Z}(I)$ .  
Primary decomposition:



- AG software ‘Macaulay 2’
- Numeric algebraic geometry methods

$$I = I_1 \cap I_2$$

$$I_1 = \langle x + y, z^2 \rangle, \quad I_2 = \langle x - y, 2y^3 - z^2 \rangle$$

$$I = I_1 \cap I_2 \cap I_3 \cap I_4 \cap I_5 \cap I_6$$

4D massless dbox hepta-cut: 6 branches of solutions

dictionary

Algebra

height  $I$

arithmetic genus

Geometry

$\dim \mathcal{Z}(I) = n - \text{height } I$  (# free parameters)

(geometric) genus

(topology)

High genus examples: arXiv:1302.1203, Rijun and YZ

works for arbitrary number of loops, any dimension

# More examples

Dimension	Diagram	# SP (ISP+RSP)	#terms in integrand basis (non-spurious + spurious)	# Solutions (dimension)
4		8 (4+4)	32 (16+16)	6 (1)
4		8 (5+3)	69 (18+51)	4 (2)
4		4 (3+1)	42 (12+30)	1(5)
4		8 (3+5)	20 (10+10)	2(2)
4		8 (4+4)	38 (19+19)	8 (1)
4- $2\epsilon$		11 (7+4)	160 (84+76)	1(4)
4		12 (7+5)	398 (199+199)	14 (2)

Three-loop!

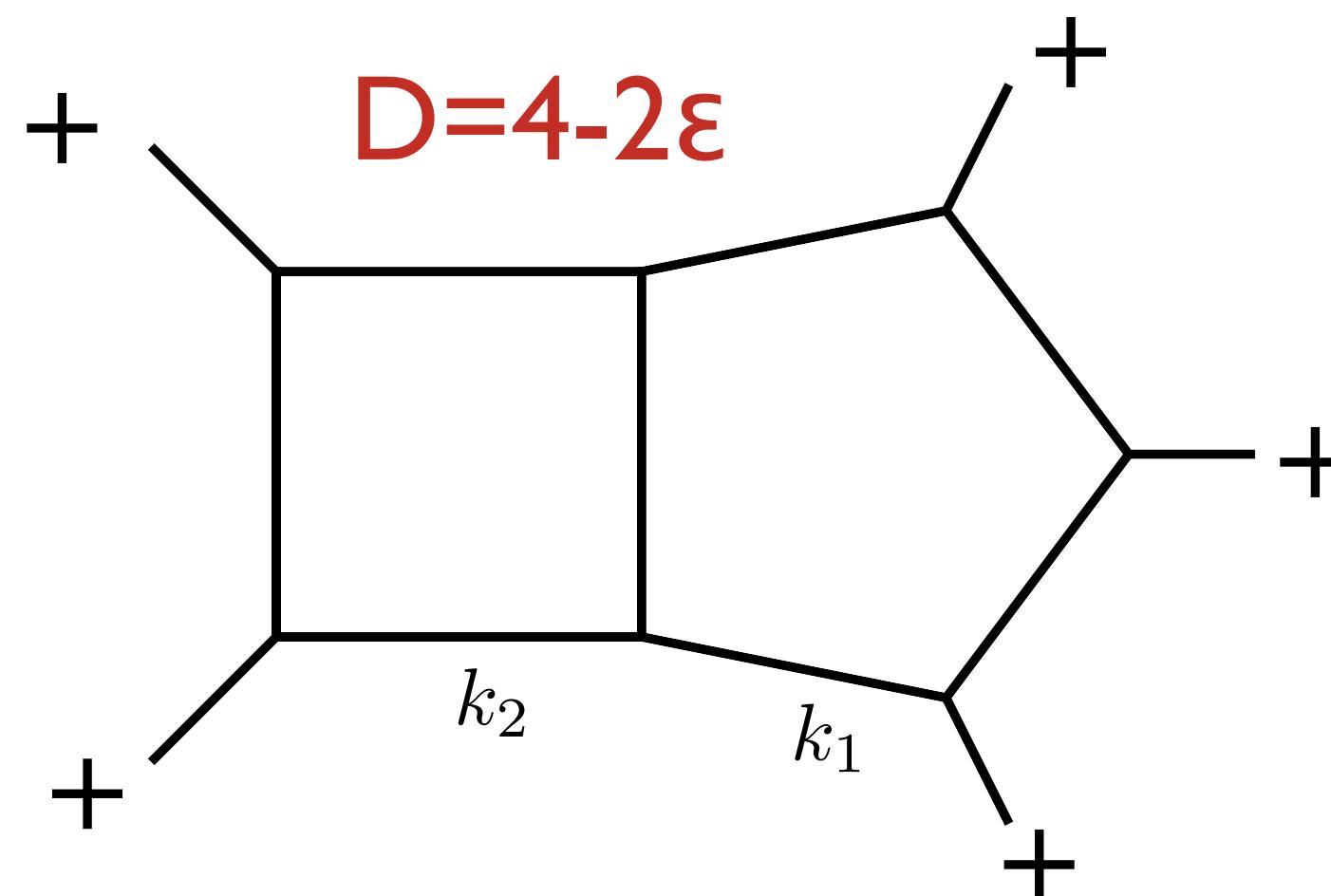
Even more examples:  
arXiv:1209.3747 Bo Feng and Rijun Huang

Nontrivial dimension  
Non-planar

# D-dim integrand reduction

2-loop 5-point QCD

arXiv: 1310.1051: Simon Badger, Hjalte Frellesvig and YZ



$$\mu_{11} = k_{[-2\epsilon],1}^2, \mu_{22} = k_{[-2\epsilon],2}^2 \text{ and } \mu_{12} = 2(k_{[-2\epsilon],1} \cdot k_{[-2\epsilon],2})$$
$$\mu_{33} = \mu_{11} + \mu_{22} + \mu_{12}$$

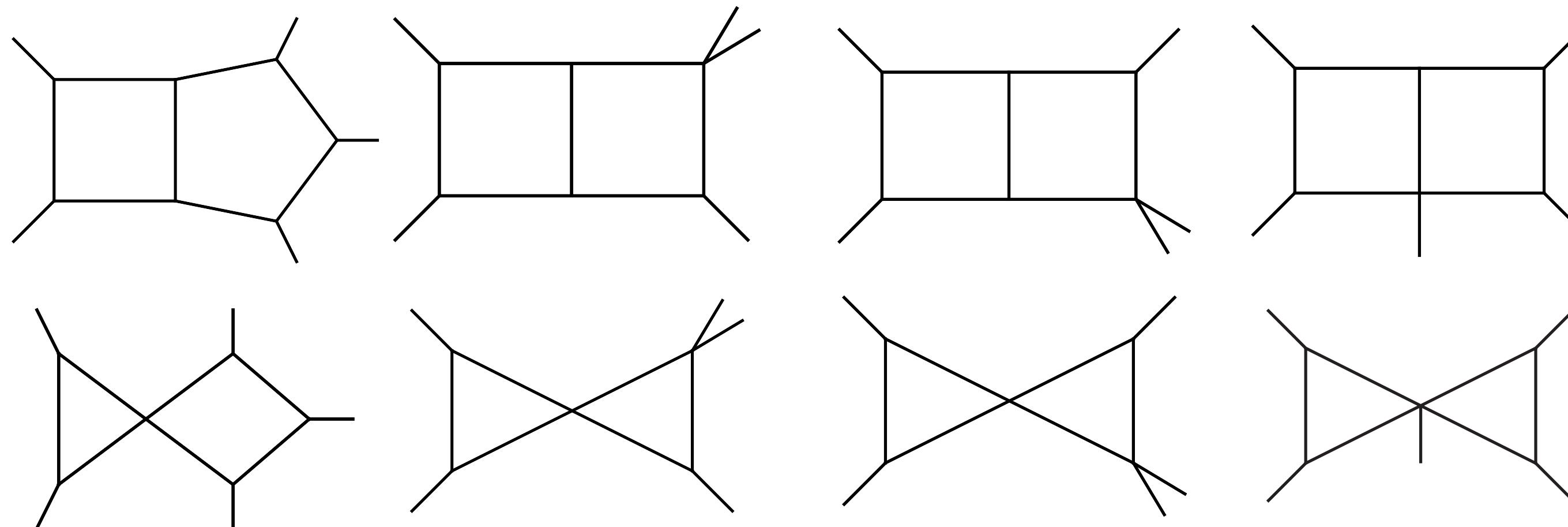
$$\Delta_{431}(1^+, 2^+, 3^+, 4^+, 5^+) = \frac{i s_{12} s_{23} s_{45} F_1(D_s, \mu_{11}, \mu_{22}, \mu_{12})}{\langle 12 \rangle \langle 23 \rangle \langle 34 \rangle \langle 45 \rangle \langle 51 \rangle} (tr_+(1345)(k_1 + p_5)^2 + s_{15}s_{34}s_{45})$$

$$F_1(D_s, \mu_{11}, \mu_{22}, \mu_{12}) = (D_s - 2)(\mu_{11}\mu_{22} + \mu_{11}\mu_{33} + \mu_{22}\mu_{33}) + 4(\mu_{12}^2 - 4\mu_{11}\mu_{22})$$

- Feynman rules + cut solution
- 6D spinor helicity formalism

Momentum-twistor parametrization  
 $(\lambda, \tilde{\lambda}) \longrightarrow (\lambda, \mu)$  (Andrew Hodges)

# 2-loop 5-gluon amplitude



arXiv: 1310.1051

first result on 2-loop 5-gluon  
helicity amplitude in QCD

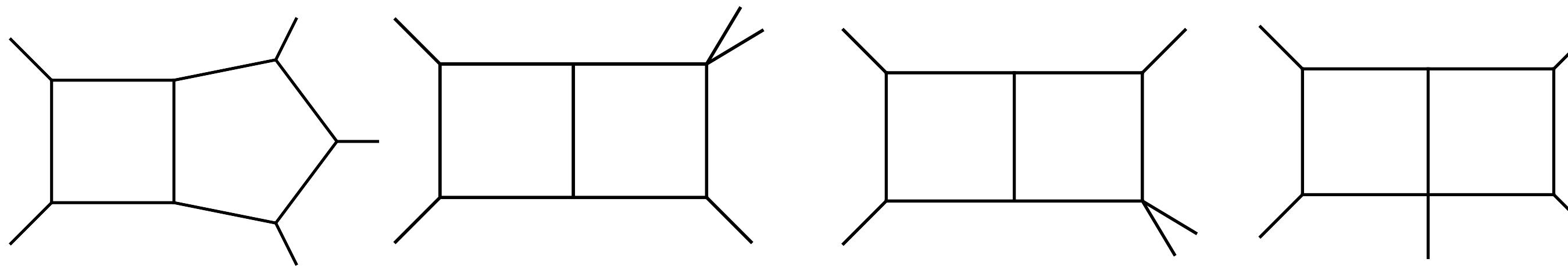
subtraction

A diagrammatic subtraction is shown. On the left, a complex loop structure with red 'x' marks on its internal lines is subtracted from a term involving the integral of a four-point function  $\Delta_{431}$ . This is followed by an arrow labeled "Integrand reduction" pointing to a simplified diagram representing the amplitude  $\Delta_{430}$ .

$$\text{Diagram with red 'x' marks} - \frac{1}{(k_1 + k_2)^2} \Delta_{431} \xrightarrow{\text{Integrand reduction}} \Delta_{430}$$

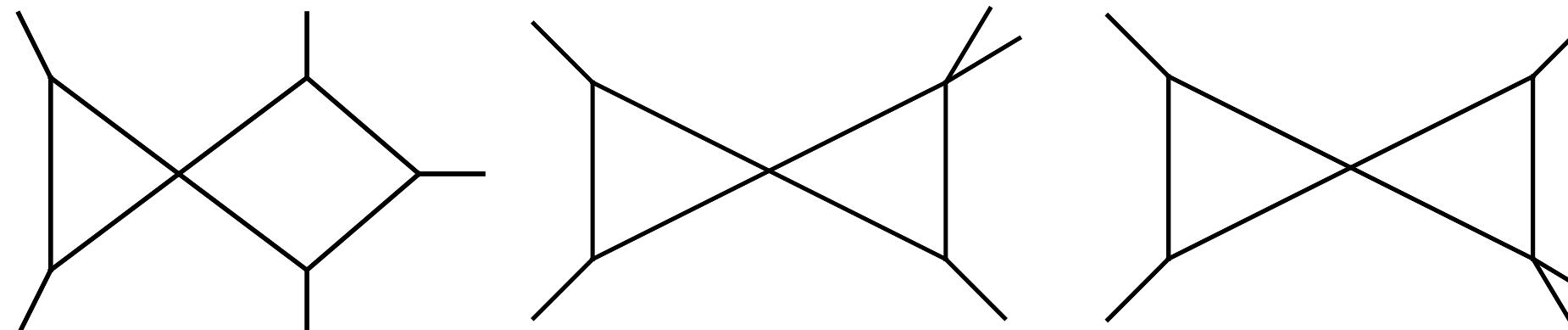
all coefficients are analytically found  
IR structure: consistent with Catani's factorization

# 2-loop 5-gluon amplitude

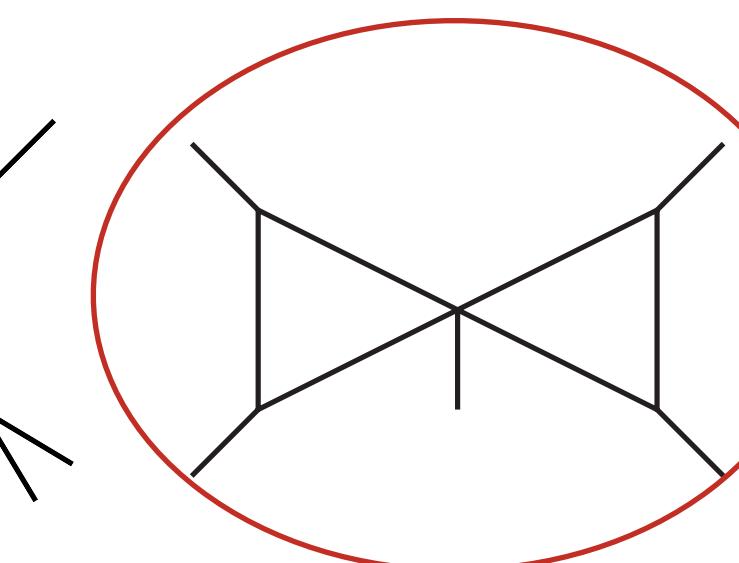


arXiv: 1310.1051

$$= F_1(D_s, \mu_{11}, \mu_{22}, \mu_{12}) \times (\text{helicity factor}) \times (\mathcal{N} = 4 \text{ Integrand})$$



No corresponding  $\mathcal{N} = 4$  diagrams



$$\begin{aligned} \Delta_{330;5L}(1^+, 2^+, 3^+, 4^+, 5^+) = & -\frac{i}{\langle 12 \rangle \langle 12 \rangle \langle 12 \rangle \langle 12 \rangle \langle 12 \rangle} \times \\ & \left( \frac{1}{2} \left( \text{tr}_+(1245) - \frac{\text{tr}_+(1345)\text{tr}_+(1235)}{s_{13}s_{35}} \right) \left( 2(D_s - 2)(\mu_{11} + \mu_{22})\mu_{12} \right. \right. \\ & + (D_s - 2)^2 \mu_{11}\mu_{22} \frac{4(k_1 \cdot p_3)(k_2 \cdot p_3) + (k_1 + k_2)^2(s_{12} + s_{45}) + s_{12}s_{45}}{s_{12}s_{45}} \Big) \\ & + (D_s - 2)^2 \mu_{11}\mu_{22} \left[ (k_1 + k_2)^2 s_{15} \right. \\ & + \text{tr}_+(1235) \left( \frac{(k_1 + k_2)^2}{2s_{35}} - \frac{k_1 \cdot p_3}{s_{12}} \left( 1 + \frac{2(k_2 \cdot \omega_{453})}{s_{35}} + \frac{s_{12} - s_{45}}{s_{35}s_{45}} (k_2 - p_5)^2 \right) \right) \\ & \left. \left. + \text{tr}_+(1345) \left( \frac{(k_1 + k_2)^2}{2s_{13}} - \frac{k_2 \cdot p_3}{s_{45}} \left( 1 + \frac{2(k_1 \cdot \omega_{123})}{s_{13}} + \frac{s_{45} - s_{12}}{s_{12}s_{13}} (k_1 - p_1)^2 \right) \right) \right] \right). \end{aligned}$$

similar for **non-planar** diagrams,  
under progress

Simon Badger, Donal O'Connell, Hjalte Frellesvig and YZ

# Momentum-twistor parametrization

Analytic computation

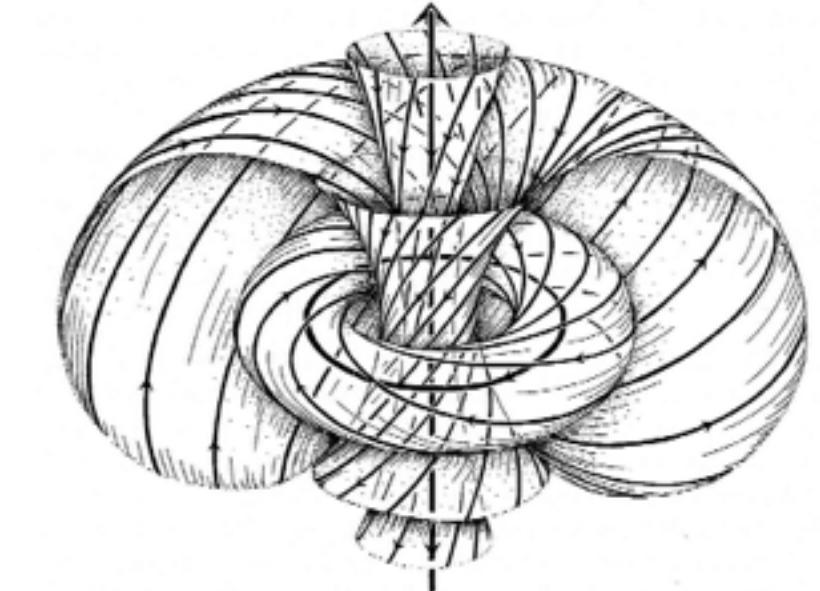
Andrew Hedges

Spinor helicity formalism  $(\lambda, \tilde{\lambda}) \longrightarrow$  Momentum-twistor parametrization  $(\lambda, \mu)$

- momentum conservation
- Schouten identity
- Fierz identity
- ...

all constraints resolved

$$\tilde{\lambda}_i = \frac{\langle i, i+1 \rangle \mu_{i-1} + \langle i+1, i-1 \rangle \mu_i + \langle i-1, i \rangle \mu_{i+1}}{\langle i, i+1 \rangle \langle i-1, i \rangle}$$



5-point

$$\begin{pmatrix} \lambda_1 & \lambda_2 & \lambda_3 & \lambda_4 & \lambda_5 \\ \mu_1 & \mu_2 & \mu_3 & \mu_4 & \mu_5 \end{pmatrix} = \begin{pmatrix} 1 & 0 & \frac{1}{x_1} & \frac{1}{x_1} + \frac{1}{x_2} & \frac{1}{x_1} + \frac{1}{x_2} + \frac{1}{x_3} \\ 0 & 1 & 1 & 1 & 1 \\ 0 & 0 & 0 & x_4 & 1 \\ 0 & 0 & 1 & 1 & \frac{x_5}{x_4} \end{pmatrix}$$

In the final result, it is easy to convert  $\{x_1, x_2, x_3, x_4, x_5\}$  to  $s_{ij}, tr_5\dots$

n-point, under progress

# Conclusion

- Algebraic geometry approach to high-loop amplitudes
  - Gröbner Basis → Integrand basis
  - Primary decomposition → Global unitarity cut structure
  - Multivariate residues → maximal unitarity

# Outlook

- Global residues via algebraic geometry  
works by K. Larsen, D. Kosower, M. Sogaard, YZ
- Integration-by-parts identities from the viewpoint of differential geometry

YZ 1406.xxxx, to appear